

BRIEF COMMUNICATION

THE EINSTEIN VISCOSITY CORRECTION IN n DIMENSIONS

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Abstract—The effective viscosity of a dilute suspension of rigid n -dimensional hyperspheres in a viscous fluid at small particle Reynolds numbers is determined; the result being

$$\mu_{\text{eff}} = \mu \left(1 + \frac{n+2}{2} \phi \right).$$

Expressions are also given for the n -dimensional Stokes velocity and pressure fields for a hypersphere in a pure straining flow.

INTRODUCTION

In teaching suspension mechanics as a topic in a course on transport processes, one is confronted with the difficulty of assigning a homework problem which is doable in a time scale small compared with that of a normal doctoral thesis. After having presented Einstein's correction to the viscosity of a dilute suspension of spheres, the natural homework problem is that for 2-dimensional rigid cylinders (a solution to Stokes equations existing in this case), and the n -dimensional case provides good practice with cartesian tensorial manipulations and the summation convention. Thus, the motivation for this brief communication.

We consider a force-free, couple-free hypersphere of radius a located at the origin in a pure straining motion with rate of strain tensor E_{ij} , which is symmetric and traceless. Making use of the linearity of the Stokes equations and the fact that the only vector present is \mathbf{x} the position vector, and following Brenner (1981), we may write for the velocity and pressure fields in cartesian tensor notation with the summation convention

$$u_i(\mathbf{x}) = E_{ij}x_j + V_{ijk}E_{jk}, \quad [1]$$

$$p(\mathbf{x}) = p^\infty + P_{jk}E_{jk}. \quad [2]$$

Here $E_{ij}x_j$ and p^∞ are the undisturbed velocity and pressure fields at infinity, and V_{ijk} and P_{jk} , being purely geometric, are given by

$$V_{ijk} = -\frac{1}{2}(\delta_{ij}x_k + \delta_{ik}x_j) \left(\frac{a}{r}\right)^{(n+2)} - \frac{(n+2)}{2} a^n \frac{x_i x_j x_k}{r^{(n+2)}} \left[1 - \left(\frac{a}{r}\right)^2 \right], \quad [3]$$

$$P_{jk} = -\mu(n+2)a^n \frac{x_j x_k}{r^{(n+2)}}, \quad [4]$$

where $r = |\mathbf{x}|$ and μ is the fluid viscosity. Substituting $n=2$ or $n=3$, one recovers the well-known results for a cylinder or sphere. The case $n=1$ is degenerate because E_{ij} (being traceless) must be zero; there can be no straining motion in one dimension.

To compute the effective viscosity of a dilute suspension of hyperspheres we form the volume average of the stress tensor σ_{ij} (Batchelor 1970) to obtain

$$\langle \sigma_{ij} \rangle = \text{I.T.} + 2\mu \langle E_{ij} \rangle + \bar{N} \langle S_{ij} \rangle, \quad [5]$$

where $\langle \rangle$ denotes volume average, \bar{N} is the number density hyperspheres, and $\langle S_{ij} \rangle = 1/N \sum_{\alpha=1}^N S_{ij}$ is the average extra particle stress; the sum being over all particles in the volume. I.T. stands for an isotropic term which is of no importance. For rigid particles the extra particle stress S_{ij} is given by

$$S_{ij} = \int \left\{ \sigma_{ik} x_j n_k - \frac{1}{3} \delta_{ij} \sigma_{lk} x_l n_k \right\} dS, \quad [6]$$

the integration being over the hypersphere surface. A direct calculation of [6] gives

$$S_{ij} = \mu (n+2) \frac{a^n}{n} \frac{2\pi^{n/2}}{\Gamma\left(\frac{1}{2}n\right)} E_{ij}$$

where $\Gamma(x)$ is the gamma function. Whence, the effective viscosity

$$\mu_{\text{eff}} = \mu \left(1 + \frac{n+2}{2} \phi \right) \quad [7]$$

where

$$\phi = \bar{N} \frac{a^n}{n} \frac{2\pi^{n/2}}{\Gamma\left(\frac{1}{2}n\right)}$$

is the volume fraction of hyperspheres. $n = 3$ gives the Einstein correction of $5/2\phi$ (Einstein 1906), and $n = 2$ gives the result for rigid cylinders first reported by Belzons *et al.* (1981).

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